

Central extensions of loop groups and loop algebras

§1

G : compact connected Lie group.

If G commutative, then $G \cong \mathbb{C}_1^r$ ($\mathbb{C}_1 = \{z \in \mathbb{C} : |z|=1\}$) a torus

So in general $Z(G)^\circ$ (id. comp. of center) a torus

$G' := (G, G)$ connected compact with finite center. Then $G = Z(G)^\circ \cdot G'$, $Z(G')$ finite.

Assume from now on $Z(G)$ finite ($\Leftrightarrow G = G'$).

Facts: Then $\pi_1 G$ finite, $\pi_2 G = \{1\}$, $\pi_3 G$ free abelian, of rank r , say.

So G has a finite universal cover $\tilde{G} \rightarrow G$ which is 2-connected and $\pi_3 \tilde{G} \cong \mathbb{Z}^r$.

A connected Lie group is called simple if every proper normal subgroup is finite.

Facts 1) $\tilde{G} = \tilde{G}_1 \cdot \dots \cdot \tilde{G}_r$ with each \tilde{G}_i simple and $\tilde{G}_i \cap \tilde{G}_j$ ($i \neq j$) contained in $Z(\tilde{G})$ hence finite.

2) $\tilde{G} = \tilde{G}_1 \times \dots \times \tilde{G}_r$ each \tilde{G}_i 2-connected and $\pi_3 \tilde{G}_i \cong \mathbb{Z}$.

So in a sense the simply connected, simple groups are the building blocks. Therefore:

We now assume G simply connected and simple. So G 2-connected and $\pi_3 G \cong \mathbb{Z}$.

Examples:

1) SU_n ($n \geq 2$) esp. $SU_2 \cong$ unit quaternions $\cong S^3$.

2) SO_n is simple for $n=3, n \geq 5$, but not simply connected: its π_1 is of order 2; its universal cover is a spin group (e.g., the univ. cover of $SO(3)$ is $\cong SU_2$)

3) quaternion analogue of SU_n .

[isom. classes of s.c. simple compact Lie groups \leftrightarrow root systems with connected Dynkin diagram]

The De Rham complex of G contains the left-invariant differential forms as a subcomplex with the same cohomology. This is also true if we consider left-forms that are both left and right invariant, but then d becomes zero (it is like the space of harmonic forms). This can be expressed in terms of the Lie algebra of:

$$H^*(G, \mathbb{R}) = (\wedge^* \mathfrak{g}^*)^G \leftarrow \text{acts on } \mathfrak{g}, \text{ hence on } \wedge^* \mathfrak{g}^* \text{ via adjoint action}$$

Example: let $\mathfrak{g} \times \mathfrak{g} \xrightarrow{\langle, \rangle} \mathbb{R}$ be an inner product inv under the adjoint action (exists always: G is compact and we can average over it). The infinitesimal version is invariance under $\text{ad}: \langle [X, Y], Z \rangle + \langle Y, [X, Z] \rangle = 0$, so that

$$\omega: (X, Y, Z) \in \mathfrak{g}^3 \mapsto \langle [X, Y], Z \rangle$$

is alternating and invariant under the adjoint action of G . So this gives

an element of $(\wedge^3 \mathfrak{g}^*)^G \cong H^3(G, \mathbb{R})$. Normalize it so that we get

a generator of $H^3(G; \mathbb{Z})$.

§2 M manifold $C^\infty(M, G)$ a group for pointwise multiplication.

We are going to define a remarkable group homomorphism

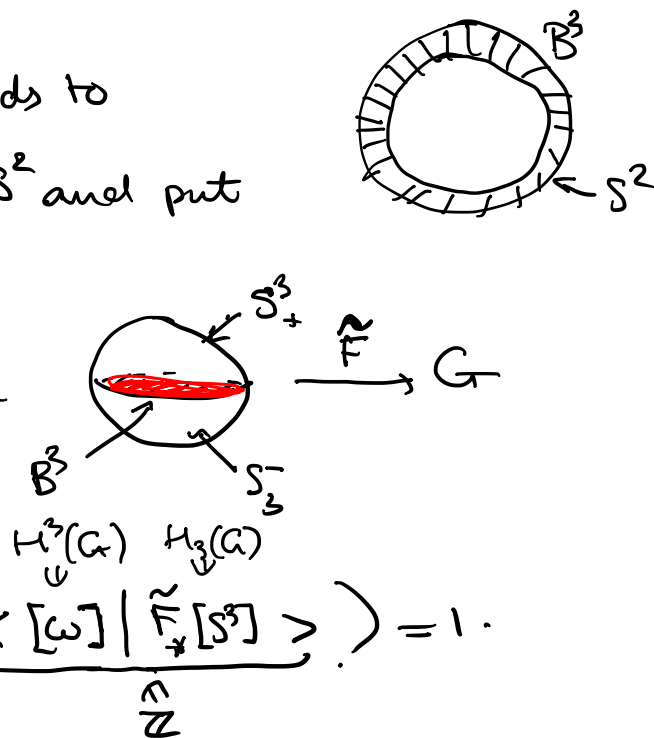
$$\Theta: C^\infty(S^2, G) \rightarrow \mathbb{C}_1$$

Let $f: S^2 \xrightarrow{C^0} G$. Since G is 2-connected, f extends to

$F: B^3 \rightarrow G$; make F constant along radii near S^2 and put

$$\Theta(F) := \exp \left(2\pi i V^{-1} \int_{B^3} F^* \omega \right) \in \mathbb{C}_1$$

Other choice F' for F combines with F to define
(F on S^3_+ ; F' on S^3_-)



$$\text{Then } \frac{\Theta(F)}{\Theta(F')} = \exp \left(2\pi i V^{-1} \int_{S^3} \tilde{F}^* \omega \right) = \exp \left(2\pi i V^{-1} \underbrace{\langle [\omega] | \tilde{F}_* [S^3] \rangle}_{\in \mathbb{Z}} \right) = 1.$$

hence $\Theta(F)$ only depends on f .

This Θ extends immediately to $C^{[\infty]}(S^2, G) := \{ \text{continuous piecewise diff maps } S^2 \rightarrow G \}$

Every element of $C^{[\infty]}(B^2, G)$ whose restriction to $\partial B^2 = S^1$ is constant 1 defines an elt. of

$C_*^{[\infty]}(S^2, G)$ (an elt which takes base point of S^2 to $1 \in G$)

So we have an exact sequence of groups

$$1 \rightarrow C_*^{[\infty]}(S^2, G) \rightarrow C^{[\infty]}(B^2, G) \rightarrow C^{[\infty]}(S^1, G) \rightarrow 1$$

You may check that $\Theta|_{C_*^{[\infty]}(S^2, G)}$ is invariant under conjugation with elements of $C^{[\infty]}(B^2, G)$. In particular, its kernel is normal in $C^{[\infty]}(B^2, G)$. If we divide out by this kernel (a "pushout") and write $\mathcal{L}G$ for $C^{[\infty]}(S^1, G)$ (referred to as the loop group of G) we get the central extension

$$1 \rightarrow \mathbb{C}_1 \rightarrow \widehat{\mathcal{L}G} \rightarrow \mathcal{L}G \rightarrow 1$$

This central extension is important because $\widehat{\mathcal{L}G}$ has a better representation theory than $\mathcal{L}G$: we have a class of irreducible reps on pre-Hilbert spaces (so-called highest weight rep's) which are in general nontrivial on \mathbb{C}_1 , the ones which are trivial on \mathbb{C}_1 factoring over G (given by evaluation at $1 \in \mathbb{C}_1$). So this gives only projective representations $\mathcal{L}G$. This is not uncommon in quantum theory, where a "state" is a point of a projectivized Hilbert space.

Its counterpart for the Lie algebra ($\mathcal{L}\mathfrak{g} := C^\infty(S', \mathfrak{g})$) is a so-called affine Kac-Moody algebra:

$$0 \longrightarrow \mathbb{R} \longrightarrow \hat{\mathcal{L}}\mathfrak{g} \longrightarrow \mathcal{L}\mathfrak{g} \longrightarrow 0,$$

$\mathcal{L}\mathfrak{g} \oplus \mathbb{R}$ as \mathbb{R} -v.m.

where the Lie bracket is given as follows: if $\alpha, \beta: S' \rightarrow \mathfrak{g}$, then

$$[\alpha + \mathbb{R}, \beta + \mathbb{R}] = ([\alpha, \beta], \frac{1}{2\pi} \int \langle \alpha(\theta), \beta(\theta) \rangle d\theta)$$

This can also be done algebraically: if k is a S' field of char zero and K a local which contains k and has k as residue field (e.g. $K = k((t))$). We assume \mathfrak{g} defined over k . We then have a central extension of $\mathfrak{g}(K) := K \otimes_k \mathfrak{g}$ by k :

$$0 \rightarrow k \rightarrow \hat{\mathfrak{g}}(K) \rightarrow \mathfrak{g}(K) \rightarrow 0$$

$$[\alpha + k, \beta + k] = ([\alpha, \beta], \text{Res} \langle \alpha, d\beta \rangle)$$

There is also an algebraic counterpart for the construction of $\hat{\mathcal{L}}G$: here G is a simple simply connected alg group / k (i.e. $G' = G$ and G has no nontrivial covers) and one can construct an extension of $G(K)$ by k by means of Steinberg symbols

§3 Let Σ now be an oriented connected closed surface, and $\overset{P}{\downarrow} \pi$ a G principal bundle (so G acts freely on the right of P , its orbits are the fibers of π and π is loc. trivial) Since G is 1-connected, such a bundle admits a section s (this trivializes the bundle) If s' is another section, then since G is 2-connected, s and s' are homotopic: if s' is given by $s'(x) = s(x)g(x)$ with $g: \Sigma \rightarrow G$, then g is homotopically trivial. If $\tilde{g}: [0,1] \rightarrow G$ is a C^∞ -homotopy (with $\tilde{g}(x,1) = g(x)$, $\tilde{g}(x,0) = 1 \in G$), then we put

$$\Theta(s) := \exp \left(2\pi\sqrt{-1} \int_{[0,1] \times G} \tilde{g}^* \omega \right) \in \mathbb{C}_1$$

and conclude as before that $\Theta(s)$ only depends on the pair (s, s') ; let us therefore denote it by $\Theta_{s'}^s$. So if we divide out by the relation $\sigma \sim \sigma' \iff \Theta_{s'}^{\sigma} = 1$, then the quotient is a \mathbb{C}_1 -torsor. This also makes sense if $\partial\Sigma \neq \emptyset$ provided we have already fixed a section (= trivialization) over $\partial\Sigma$. This is even of interest when $(\Sigma, \partial\Sigma) \in (B^2, S^1)$.

This has been used in Jones-Witten theory and (related to this) conformal field theory. One then considers a compact Riemann surface C and a simple simply conn. alg grp G over \mathbb{C} . The principal G -bundles over C have a moduli stack $\mathcal{M}(C, G)$, and the alg counterpart constructs an ample line bundle $\mathcal{L} \downarrow_{\mathcal{M}(C, G)}$ which generates $\text{Pic}(\mathcal{M}(C, G))$. The space of sections $\mathcal{L}^{\otimes l}$ is finite-dimensional - the conformal blocks of level l .